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Crystal Structures of D,L-Homocysteine Thiolactone Hydrochloride: Two Polymorphic Forms and a Hybrid*

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The structures of three different crystalline forms of D,L-homocysteine thiolactone hydrochloride have been determined by X-ray diffraction techniques. Two of the crystal forms are polymorphs: one crystallizes in space group *Pbca* and the other crystallizes in space group *Pbc2₁*. The third form is a hybrid in which both of these polymorphs coexist. The *Pbca* and *Pbc2₁* domains of the hybrid crystal scatter X-rays independently but have their corresponding unit-cell axes oriented in the same directions.

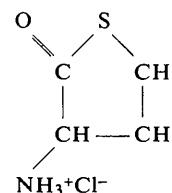
The structures of the polymorphs were arrived at by conventional heavy-atom Fourier and vector-coincidence methods, and each structure was refined by two different least-squares criteria, one based on reflection amplitudes and the other based on reflection intensities. To verify the structure of the hybrid form, a linear combination of reflection intensities from the *Pbca* and *Pbc2₁* crystals was scaled by least squares to the reflection intensities of the hybrid crystal; the agreement obtained between the reflection intensities corresponds to a conventional *R* index of 9%.

The crystal structures of the two polymorphs are very similar. The five-membered thiolactone ring is in the envelope conformation, with four atoms lying in a plane from which the β carbon atom is displaced by approximately 0.7 Å. In both structures the molecules are packed in layers perpendicular to the unit-cell *c* axis, and there is a network of hydrogen bonds between the ammonium group and chloride ions. Apparent small differences in bond-lengths and angles between molecules in the two structures are not significant; differences of about the same magnitude were produced by the two alternate methods of refinement for each individual polymorph.

Introduction

The structures of three different crystalline forms of D,L-homocysteine thiolactone hydrochloride (HTL)

* Most of the material in this communication has been abstracted from a Ph.D. thesis submitted by S.T.F. to the University of Washington in 1964.



will be described in this communication. The structures are very closely related; in fact one of them is a mixture of the other two, with at least the outward appearance of an ordinary single crystal. Crystals of all three forms were grown under similar conditions.

The first crystals we obtained happened to be the hybrid form (but we were not aware of this until much later, after the structures of the other forms had already been determined) with apparent space group *Pbca*, but pseudo *Pbc*₂. When two plausible trial structures, one in each of these space groups, failed to refine below an *R* index of 21%, we decided to recrystallize the HTL in the hope of obtaining better data crystals. Instead, a new crystal form with authentic space group symmetry *Pbc*₂ was discovered. When the new intensity data were used, it then became possible to refine the *Pbc*₂ trial structure to an acceptably low value for the *R* index.

However, renewed attempts to solve the structure of the hybrid crystal continued to fail, and so the recrystallization experiments were resumed. Eventually we obtained *Pbca* crystals which were *not* hybrids, and once again a trial structure deduced from the original data, this time the *Pbca* trial structure, was readily refined with data obtained from the new crystal form.

Finally, by comparing reflection intensities from the original crystal form with reflection intensities from the two well-behaved forms, it was possible to demonstrate that the original crystal was actually a hybrid in which the two polymorphs coexist.

The two polymorphic forms of HTL will be referred to collectively as 'polymorphs' and individually as HTL-*Pbca* and HTL-*Pbc*₂ throughout the rest of this communication. We have adopted the terminology of Ubbelohde & Woodward (1946*b*) in designating the third crystal form 'hybrid' HTL.

Experimental

All crystals were grown at room temperature from solutions in methanol-nonsolvent mixtures. The hybrid and *Pbc*₂ crystals were obtained with HTL purchased from Mann Research Laboratories, Inc., and the *Pbca* crystals with HTL from the California Corporation for Biochemical Research.

The hybrid crystal form, HTL-hybrid, was crystallized by slow evaporation from a solution of HTL in a mixture of methanol and amyl alcohol. The crystals exhibited a lath-like habit with prominent (100) faces and with the long dimension parallel to the unit-cell *c* axis. Although most crystals large enough for data collection were severely cracked, there was no evidence of multiple growth. Systematic absence of reflections

*Ok**l* with *k* odd, *h*0*l* with *l* odd, and *hk*0 with *h* odd apparently established the space group as *Pbca*. However, the extreme weakness of most *h* odd reflections indicated pseudo-*Pbc*₂ symmetry (with the *a* axis halved). All reflections were sharp, and there was no indication of streaking. The data crystal, a fragment approximately 0.1 mm on a side, was dipped repeatedly into liquid nitrogen in an attempt to reduce possible extinction effects.

HTL-*Pbc*₂ crystallized from a methanol solution in vapor equilibrium with methanol-dioxan (1:1). The crystal habit was similar to that of HTL-hybrid. The crystal density and systematic absence of reflections *Ok**l* for *k* odd and *h*0*l* for *l* odd established its space group as *Pbc*₂ (space group no. 29, *Pca*₂, for the setting given in *International Tables*). The dimensions of the data crystal were 0.08 × 0.1 × 0.12 mm.

Well-formed large plates, (100) faces prominent, of HTL-*Pbca* crystallized from an evaporating solution in methanol and amyl alcohol. The same systematic absences were observed as for HTL-hybrid, establishing the space group as *Pbca*. Data were collected with a crystal fragment approximately 0.1 mm along each edge.

Unit-cell constants for the three crystal forms are given in Table 1. Lattice parameters for HTL-hybrid were measured by precession photography; uncertainties were estimated from repeated measurement of the precession films. Lattice parameters and their estimated standard deviations for HTL-*Pbca* and HTL-*Pbc*₂ were determined by least-squares treatment of 2θ angles measured with a General Electric XRD-5 diffractometer. Copper Kα radiation was used throughout, for both intensity and lattice parameter data, a wavelength of 1.5418 Å being assumed for the mean of Kα₁ and Kα₂, or 1.54051 Å for Kα₁ where the two components were resolvable in high order reflections; all measurements were made at room temperature. Densities were obtained by free flotation in carbon tetrachloride-bromobenzene. The uncertainties are r.m.s. deviations among six independent measurements for HTL-hybrid and three independent measurements for each of HTL-*Pbc*₂ and HTL-*Pbca*.

Intensities were collected on the diffractometer by a 2θ-scanning technique; stationary background counts were taken at both the start and finish of the peak scan. All unique reflections within range of the instrument (2θ ≤ 160°) were examined; 1452 reflections measured, with 66% observed for HTL-hybrid; 796 reflections measured, with 86% observed for HTL-*Pbc*₂; and 1413 reflections measured, with 90% observed for HTL-*Pbca*. For HTL-hybrid the criteria of observability were left to the judgment of the diffractometer

Table 1. Unit-cell constants

	<i>a</i>	<i>b</i>	<i>c</i>	<i>D</i> _{meas}	<i>D</i> _{calc}	<i>Z</i>
HTL-Hybrid	19.55 ± 0.03 Å	9.29 ± 0.02 Å	7.23 ± 0.02 Å	1.535 ± 0.003 g.cm ⁻³	—	—
HTL- <i>Pbc</i> ₂	9.806 ± 0.003	9.321 ± 0.002	7.321 ± 0.002	1.519 ± 0.001	1.525 g.cm ⁻³	4
HTL- <i>Pbca</i>	19.512 ± 0.002	9.296 ± 0.003	7.272 ± 0.001	1.543 ± 0.002	1.547	8

operator and intensities of unobserved reflections were not recorded; an estimated minimum observable F_o was assigned to unobserved reflections during data reduction. For HTL- $Pbc2_1$ and HTL- $Pbca$, on the other hand, the intensities of all reflections were recorded and the criterion of observability was based on counting statistics.

Although absorption effects are significant for a compound with the elementary composition of HTL, no correction was applied, and this is probably a source of appreciable error. Nevertheless, since all the intensity data were collected with small and approximately equidimensional crystal fragments, the error is probably not excessive.

Structure determination and initial refinement of HTL- $Pbca$ and HTL- $Pbc2_1$

A combination of vector-coincidence and heavy-atom Fourier methods was used to deduce a plausible trial structure in space group $Pbca$ from the HTL-hybrid intensity data. However, full-matrix least-squares refinement with a slightly modified version of the Busing & Levy (1959) program failed to reduce the R index below 21%. Since only 38% of the h -odd reflections for HTL-hybrid were observed, it was decided to begin again, ignoring the h -odd reflections entirely, and attempt to solve the structure in space group $Pbc2_1$. A second trial structure, very closely related to the first, was readily obtained, but once more it could not be refined below an R index of 21%.

Since neither trial structure could be refined with the HTL-hybrid intensity data, it seemed possible that the original data crystal was disordered in a manner somewhat analogous to azulene (Robertson, Shearer, Sim & Watson, 1962). A difference-Fourier synthesis phased on the $Pbca$ trial structure apparently verified this conclusion: the difference map contained regions of negative electron density surrounding the positions of several of the input atoms and positive peaks in locations related to these positions by reflection through a mirror at $z = \frac{1}{2}$. For this reason we suspected that HTL-hybrid was a disordered $Pbca$ structure in which alternate molecules were randomly replaced by their enantiomorph. However, trial structures which included the original $Pbca$ trial structure, plus its mirror image, as fractional atoms could not be refined. In order to eliminate the possibility that HTL-mixed contained molecules with different conformations of the thiolactone ring, we also tried to refine several trial structures in which only certain atoms were replaced by two fractional atoms. Again, no refinable trial structure was obtained. Finally, in desperation, we decided to recrystallize the HTL and start all over with new data crystals.

The subsequent discovery of two additional crystal forms, one with genuine $Pbca$ symmetry and the other with genuine $Pbc2_1$ symmetry, has already been outlined. After these crystals were obtained, we were able to refine each of the above trial structures with newly

collected intensity data from the corresponding polymorph. Atomic scattering factors given in *International Tables for X-ray Crystallography* (1962) were used throughout the calculations. For the refinement of both polymorphs a weighting scheme similar to Hughes's (1941) was applied, the weight assigned to each reflection being the smaller of the two quantities $[\sigma(F)]^{-2}$ or $(0.05F)^{-2}$, where $\sigma(F)$ was estimated from counting statistics. The quantity minimized was $\sum w(F_o - kF_c)^2$. Unobserved reflections for which F_c was less than F_o were omitted from the matrix of normal equations.

Refinement calculations of HTL- $Pbc2_1$ were carried out on an IBM-709 computer with the full-matrix least-squares program of Busing & Levy (1959). Seven cycles reduced the R index to 10.8%, at which point difference-Fourier synthesis and bond-length calculations indicated an erroneous placement of C(4). Refinement was then continued, with new C(4) coordinates, to an R index of 6.2%. The eight hydrogen atoms of the asymmetric unit were located in a new difference-Fourier synthesis, albeit with some uncertainty, and introduced into the structure with the fixed anisotropic temperature factors of the atoms to which they were covalently bonded. Three final refinement cycles were calculated in which all parameters except hydrogen-atom temperature factors were varied. The last parameter shift was in no case greater than one standard deviation. The final R index was 4.9% for 682 equally weighted observed reflections. The estimated standard deviation for an observation of unit weight, $[\sum w(\Delta F)^2/(m-s)]^{1/2}$, was 1.28.

In the case of HTL- $Pbca$, refinement calculations were carried out on a CDC-1604 machine with a modified block-diagonal version of the least-squares program of Gantzel, Sparks & Trueblood (1961). Seven cycles of refinement reduced the R index to 6.5%. The eight highest peaks (0.55–0.86 e.Å⁻³) in a difference-Fourier synthesis calculated at this point were in reasonable locations for hydrogen atoms. When these were introduced into the structure, each with the fixed anisotropic temperature factor of the atom to which it was covalently bonded, the R index dropped to 4.9%. Nevertheless the parameter shifts on the next cycle were still unreasonably large and further refinement cycles failed to cause convergence. Indeed, during subsequent cycles it became obvious that most of the parameters were oscillating with increasing amplitude about fairly well-defined mean values. Finally each parameter was assigned this mean value and a damping factor of 0.5 was applied to all subsequently calculated shifts. This procedure succeeded in causing convergence in three additional cycles. The mean final shift was 0.2σ , and no shift was larger than σ . The R index for 1265 equally weighted observed reflections was 4.2% and the estimated standard deviation for an observation of unit weight was 1.46. At a later date, when a *full-matrix* least-squares program for the CDC-1604 computer became available, it was used to calculate one additional refinement cycle in order to con-

firm the validity of the block-diagonal treatment. This cycle produced a maximum parameter shift of 0.2 σ , and a mean shift of only 0.08 σ .

Difference-Fourier syntheses, phased on final parameters for all atoms, were calculated for both polymorphs as a concluding check on the refinement. No interpretable features could be seen.

Deduction of the crystal structure of HTL-hybrid

Even though essentially correct trial structures for both polymorphs had been derived from the intensity data of HTL-hybrid, the structure of this crystal form was as yet unknown. Two possible explanations seemed reasonable: (1) HTL-hybrid was disordered in a way which had not occurred to us, and (2) HTL-hybrid was composed of separate domains of HTL-*Pbca* and HTL-*Pbc2₁* which scattered X-rays independently but had their unit-cell axes oriented in the same directions. Since several composite trial structures (the calculated structure factors for one of which were equivalent to a linear combination of F_o^2 's of the two polymorphs) designed to allow for certain kinds of random disorder in space group *Pbca* had already failed to refine, we decided to test the second of these hypotheses. The idea that there were independent scattering regions within the HTL-hybrid crystal was further supported when high-order reflections for three freshly grown HTL-hybrid single crystals were found to exhibit pronounced splitting. For example, four distinct peaks, with 2θ values of 147.7°, 147.9°, 149.1°, and 150.2°, were observed for the 24,2,0 reflection. These values correspond reasonably well with Bragg angles calculated for Cu $K\alpha_1$ and α_2 from lattice parameters of HTL-*Pbc2₁* (146.2° and 147.2°) and HTL-*Pbca* (148.2° and 149.2°).

As a further test of the hybrid-crystal hypothesis, a linear combination of F_o^2 (*Pbca*) and F_o^2 (*Pbc2₁*), each on an absolute scale, was scaled by least-squares to F_o^2 (hybrid), and the agreement index of the resulting 'calculated intensities' for the best linear combination was compared with the corresponding agreement obtained from the intensities for either polymorph alone. Specifically, scale constants k_1 and k_2 were adjusted to minimize $\Sigma\{F_o^2(\text{hybrid}) - [k_1F_o^2(\text{Pbca}) + k_2F_o^2(\text{Pbc2}_1)]\}^2$, but were otherwise unconstrained. The agreement index was calculated as $R_2 = \Sigma|F_o^2(\text{hybrid}) - [k_1F_o^2(\text{Pbca}) + k_2F_o^2(\text{Pbc2}_1)]| / \Sigma F_o^2(\text{hybrid})$. Only 863 reflections which were sufficiently strong to be considered 'observed' for all crystal forms were included in the summations.

Results were as follows:

	R_2	k_1	k_2
HTL (<i>Pbca</i>) only	36.7%	5.15	0
HTL (<i>Pbc2₁</i>) only	25.5%	0	23.42
Linear combination	11.2%	2.12	15.10

Comparable results were also obtained when the F_o^2 instead of F_o of the polymorphs were similarly scaled

to the F_o^2 of the hybrid crystal; corresponding R_2 values were 35.7% for *Pbca* only, 27.0% for *Pbc2₁* only, and 13.1% for the best linear combination. The conventional R index, taking F_c (hybrid) to be $[k_1F_o^2(\text{Pbca}) + k_2F_o^2(\text{Pbc2}_1)]^{1/2}$, was 8.6% for the 863 'observed' reflections.

We believe that the much smaller value for R_2 obtained with the best linear combination, when compared with the value for R_2 obtained with either polymorph alone, is fairly good verification of the hybrid structure proposed for HTL-hybrid.

Three estimates of the relative number of HTL molecules in *Pbca* and *Pbc2₁* regions of the hybrid crystal were made. These were based upon (1) the observed densities of the three crystal forms; (2) the least-squares scale constants, k_1 and k_2 ; and (3) comparison of the ratio $P(\frac{1}{2}, 0, 0)/P(0, 0, 0)$ in the HTL-hybrid and HTL-*Pbca* Patterson maps (the origin peak from the HTL-*Pbc2₁* portion contributes to $P(\frac{1}{2}, 0, 0)$ in a Patterson map calculated with the HTL-hybrid intensities). The values obtained, expressed as per cent of the molecules in HTL-*Pbc2₁* regions, are 32%, 64%, and 58%, respectively. Although the estimates vary considerably, it should be noted that the low figure based upon density measurements is probably the most inaccurate. In any case, all three results do indicate that the hybrid crystal form must be composed of roughly equal amounts of the constituent polymorphs.

The arrangement of *Pbca* and *Pbc2₁* regions within the hybrid crystal remains unknown. As far as our data are concerned, the size of individual *Pbca* and *Pbc2₁* regions could range from separate randomly intermixed mosaic blocks, on the one hand, to regions so large, on the other hand, that the crystal would be considered a heterogenous twin.

The existence of different subcrystalline domains within a 'single' crystal was first observed by Ubbelohde & Woodward while studying continuous phase transitions of Rochelle salt (1946a) and potassium dihydrogen phosphate (1946b). Since then, crystals of many other compounds which undergo continuous phase transitions have also been observed to change into crystal hybrids near the transition temperature (Ubbelohde, 1963). It would be of interest to know whether these polymorphs of HTL could exist in equilibrium at room temperature, and whether HTL-hybrid could result from a continuous phase transition. We have no data which would bear on these questions and it is possible that HTL-hybrid may instead result from co-crystallization of the two polymorphs.

F^2 Refinement of HTL-*Pbca* and HTL-*Pbc2₁*

When a full-matrix least-squares program for the CDC 1604 machine became available to us [adapted by Dr R. A. Alden from the program of Busing, Martin & Levy (1962)], we decided to try F^2 full-matrix refinement of HTL-*Pbca* and HTL-*Pbc2₁*. The quantity minimized in this refinement was $\Sigma w(F_o^2 - kF^2)^2$, where $w =$

$[1/\sigma(F_o^2)]^2$ with $\sigma(F_o^2)$ determined from counting statistics. This is to be contrasted with the quantity minimized in the refinement procedure previously described, namely $\sum w(F_o - kF_c)^2$. No distinction was now made between 'observed' and 'unobserved' reflections. Weak reflections for which background intensity accidentally exceeded peak intensity were assigned $F_o^2 = 0$ at data reduction time. All reflections except 002 contributed to the matrix of normal equations; the 002 reflection of HTL-*Pbca* exhibited obvious extinction effects and was therefore not included in either refinement or *R* index calculations. Parameters from the final least-squares cycle of *F* refinement were used as input to the F^2 refinement of both structures.

The net result of this computational experiment was that the F^2 refinement of both HTL-*Pbca* and HTL-*Pbc2*₁ converged to structures which differed slightly, but significantly, from the structures obtained with the corresponding *F* refinements. Details of the procedure are outlined in the following paragraphs.

F^2 refinement of HTL-*Pbca* progressed routinely. Two initial cycles were calculated with hydrogen thermal parameters fixed. Most of the refinement was accomplished during these cycles. Hydrogen atoms were then assigned individual isotropic temperature factors and two more cycles were run in which only hydrogen thermal parameters were permitted to vary. Finally, all parameters, both hydrogen and non-hydrogen, were adjusted in two concluding cycles. In the last cycle the maximum parameter shift was 0.28σ and the mean shift was 0.05σ . The estimated standard deviation for an observation of unit weight, $[\sum w(\Delta F^2)^2/(m-s)]^{1/2}$ was 3.03.

F^2 refinement of HTL-*Pbc2*₁ did not proceed as smoothly since the hydrogen atoms of this structure were not as well defined. Hydrogen thermal parameters became negative when they were permitted to vary. After several unsuccessful attempts at refinement, individual isotropic temperature factors were assigned to hydrogen atoms (3.1 \AA^2 for hydrogen atoms covalently bonded to carbon atoms and 4.8 \AA^2 for hydrogen atoms covalently bonded to the nitrogen atom) and were not adjusted during further refinement. Most of the refinement took place during the first two cycles, in which hydrogen positional parameters were also fixed. Finally, four concluding cycles were run, in which all parameters except hydrogen thermal parameters were permitted to vary. It is doubtful that hydrogen positional parameters actually refined significantly since they were observed to shift erratically throughout the refinement. In the last cycle the mean parameter shift was 0.2σ and the maximum shift was 1.3σ (only one parameter, $H(2)-y$, was altered by more than one standard deviation). The estimated standard deviation for an observation of unit weight was 3.54.

Over the course of the F^2 refinement of each polymorph the weighted sum of squares of the F^2 residuals decreased by approximately 50%; and, as would be

expected, parameters from the final F^2 cycle produced an *increase* in the weighted sum of squares of the residuals calculated for the *F*'s.

Differences between bond lengths and angles resulting from the two types of refinement are, on the average, of the same magnitude as the estimated standard deviations. Although the mean bond-length difference between the *F* and F^2 refinement was only 1.2σ (the mean σ is 0.003 \AA for HTL-*Pbca* and 0.01 \AA for HTL-*Pbc2*₁) for both crystal forms, the important point to be made is that *some* of the refinement-dependent differences were very appreciable, *i.e.* 0.01 \AA for HTL-*Pbca*, and 0.04 \AA for HTL-*Pbc2*₁. Because of this, no significance can safely be attributed to the small differences apparently observed between molecules in the HTL-*Pbca* and HTL-*Pbc2*₁ structures, even though some of these exceeded 2σ (see Tables 8 and 9). Possibly, elimination of systematic errors in the data, in particular application of absorption corrections, would reduce this effect.

Since it seems preferable from a theoretical point of view to base parameter refinement on quantities proportional to the first power of the data rather than the square root of the data, parameters and calculated structure factors reported for HTL-*Pbca* and HTL-*Pbc2*₁ in this communication are taken from the F^2 refinement. Final *R*₂ indices, based on F^2 and calculated over all reflections, were 7.4% for HTL-*Pbc2*₁ and 5.0% for HTL-*Pbca*; corresponding *R* indices, based on the *F*'s, were 6.5% and 4.8%.

Refined positional parameters and their estimated standard deviations are listed in Table 2. Thermal parameters are given in Table 3. Observed reflection amplitudes and calculated structure factors for the two polymorphs are listed in Tables 4 and 5. The corresponding table for the hybrid crystal structure has been omitted since presumably it represents only one of a series of hybrid structures containing varying amounts of the two components.

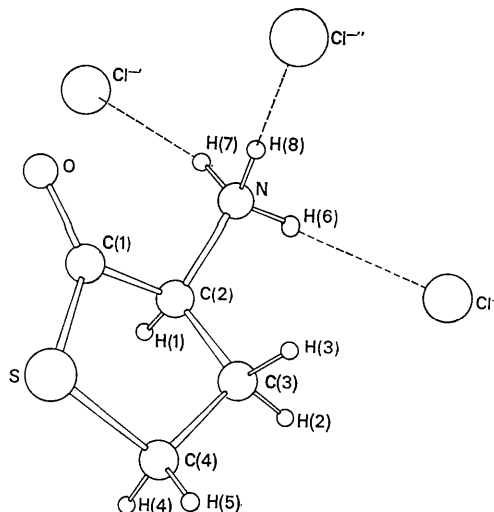


Fig. 1. Computed perspective projection of a single HTL molecule.

Results and discussion

Fig. 1 is a representation of a single HTL molecule from the HTL-*Pbca* crystal structure. The drawing is a perspective projection from an arbitrary vantage point, chosen to show the thiolactone ring and the arrangement of hydrogen bonds most clearly. Molecular packing of both polymorphs is shown in Figs. 2, 3, and 4. Hydrogen atoms are omitted for clarity. Figs. 2 and 3 are orthographic projections along the *c* axis of HTL-*Pbca* and HTL-*Pbc2*₁ respectively; hydrogen bonds are indicated by dashed lines. Orthographic projections along the *b* axis are shown in Fig. 4 for both crystal forms.

Bond lengths and angles of covalent bonds not involving hydrogen atoms are given in Tables 6 and 7. Covalent bond lengths and angles involving hydrogen atoms are listed in Table 8 for HTL-*Pbca*. Since the

hydrogen atoms of the *Pbc2*₁ structure were not well-defined, the corresponding table for this polymorph is omitted; the range of covalent bond lengths to hydrogen atoms for HTL-*Pbc2*₁ was from 0.65 to 1.42 Å. Hydrogen-bond lengths and angles are listed in Tables 9 and 10. Bond lengths and angles involving HTL-*Pbc2*₁ hydrogen atoms are again omitted.

The HTL molecules in the *Pbca* and *Pbc2*₁ crystal structures are nearly identical. Although examination of bond lengths and angles for the two structures does reveal differences, it should be emphasized again that similar differences were obtained when either polymorph was refined by two different methods. Thus, differences in bond lengths and angles of as much as 0.03 Å and 3.5° in these two structures may well be without significance.

The overall shape of the molecules is that of a puckered five-membered ring, four atoms lying nearly

Table 2. Positional parameters, in fractional coordinates $\times 10^4$ ($\times 10^3$ for hydrogen parameters), and their estimated standard deviations

Atom	HTL- <i>Pbca</i>			HTL- <i>Pbc2</i> ₁		
	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>
Cl ⁻	4066 (1)	760 (1)	2615 (1)	8192 (2)	839 (2)	2615 (*)
S	3114 (1)	6973 (1)	2246 (1)	6160 (2)	6969 (2)	2053 (5)
O	4380 (1)	6970 (2)	0963 (3)	8700 (5)	7020 (5)	0877 (9)
N	4553 (1)	4008 (2)	1786 (4)	9100 (6)	4102 (7)	1743 (12)
C(1)	3950 (1)	6347 (3)	1843 (4)	7842 (7)	6419 (7)	1763 (12)
C(2)	4041 (1)	4889 (3)	2794 (3)	8063 (7)	4953 (6)	2750 (14)
C(3)	3340 (1)	4193 (3)	2911 (4)	6671 (9)	4239 (8)	2778 (17)
C(4)	2855 (1)	5382 (3)	3512 (4)	5633 (9)	5394 (8)	3395 (15)
H(1)	417 (1)	511 (3)	399 (3)	822 (7)	552 (7)	405 (10)
H(2)	331 (1)	333 (3)	371 (4)	652 (10)	384 (9)	348 (12)
H(3)	318 (1)	377 (3)	172 (4)	652 (7)	335 (7)	122 (11)
H(4)	290 (1)	562 (3)	487 (4)	581 (6)	562 (7)	495 (11)
H(5)	237 (1)	517 (3)	307 (4)	454 (6)	517 (7)	292 (12)
H(6)	458 (1)	304 (3)	231 (5)	943 (7)	331 (6)	268 (14)
H(7)	504 (1)	431 (3)	201 (4)	986 (8)	445 (8)	191 (14)
H(8)	451 (2)	398 (3)	059 (5)	936 (8)	425 (9)	034 (12)

Table 3. Thermal parameters of non-hydrogen atoms and their estimated standard deviations

β as given here is defined by: $T = \exp[-10^{-4}(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$

	HTL- <i>Pbca</i>					
	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Cl ⁻	18 (1)	65 (1)	138 (1)	2 (1)	-6 (1)	1 (1)
S	21 (1)	79 (1)	225 (2)	16 (1)	17 (1)	31 (1)
O	22 (1)	65 (2)	208 (5)	0 (1)	16 (1)	23 (3)
N	16 (1)	56 (2)	139 (5)	1 (1)	4 (1)	10 (3)
C(1)	18 (1)	54 (3)	136 (5)	3 (1)	3 (2)	-2 (3)
C(2)	14 (1)	61 (3)	103 (5)	1 (1)	1 (2)	-3 (3)
C(3)	13 (1)	72 (3)	175 (7)	-5 (1)	5 (2)	-2 (4)
C(4)	17 (1)	82 (3)	171 (7)	-1 (1)	11 (2)	1 (4)
	HTL- <i>Pbc2</i> ₁					
	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Cl ⁻	109 (2)	99 (2)	223 (4)	5 (2)	1 (3)	-3 (3)
S	131 (2)	122 (2)	304 (7)	35 (2)	30 (4)	38 (4)
O	134 (7)	99 (7)	309 (17)	-16 (5)	13 (9)	18 (10)
N	95 (7)	91 (7)	221 (15)	4 (6)	21 (9)	20 (11)
C(1)	117 (9)	94 (8)	213 (21)	7 (7)	1 (12)	-9 (12)
C(2)	87 (7)	92 (7)	202 (17)	-5 (6)	27 (13)	-6 (13)
C(3)	98 (7)	104 (9)	294 (24)	0 (8)	33 (15)	5 (15)
C(4)	127 (11)	104 (10)	298 (25)	2 (8)	61 (15)	0 (14)

Table 4. Observed and calculated structure factors for HTL-Pbca

Within each group of constant k and l, the columns contain from left to right: h, 10F_o, and 10F_c. An asterisk designates reflections for which the observed intensity was less than 1σ above background. F_o's followed by E were affected by secondary extinction.

Table with multiple columns for h, 10F_o, and 10F_c values across various k and l indices. The table is organized into groups for different k and l values, with h values ranging from 0 to 20 and 10F values ranging from approximately -100 to 500. Asterisks indicate reflections where observed intensity was less than 1σ above background.

Table 5. *Observed and calculated structure factors for HTL-Pbc₂*

Within each group of constant *k* and *l*, the columns contain from left to right: *h*, $10F_o$, $|10F_c|$, and the phase angle in millicycles. An asterisk designates reflections for which the observed intensity was less than 1σ above background.

(Detailed grid of structure factor data for HTL-Pbc₂, including columns for h, 10F_o, |10F_c|, and phase angle in millicycles. The table is organized into multiple groups based on k and l values.)

Table 6. *Bond lengths of covalent bonds, and their estimated standard deviations*

Bond	Length	
	HTL-Pbca	HTL-Pbc ₂
S-C(1)	1.756 ± 0.003 Å	1.740 ± 0.007 Å
S-C(4)	1.814 ± 0.003	1.840 ± 0.009
O-C(1)	1.204 ± 0.003	1.201 ± 0.009
N-C(2)	1.486 ± 0.003	1.485 ± 0.010
C(1)-C(2)	1.532 ± 0.003	1.561 ± 0.010
C(2)-C(3)	1.515 ± 0.004	1.519 ± 0.011
C(3)-C(4)	1.520 ± 0.004	1.549 ± 0.012

Table 7. *Bond angles for covalent bonds, and their estimated standard deviations*

Angle	HTL-Pbca	HTL-Pbc ₂
	C(1)-S-C(4)	94.3 ± 0.1°
O-C(1)-S	125.3 ± 0.2	126.3 ± 0.5
O-C(1)-C(2)	125.8 ± 0.2	124.1 ± 0.6
S-C(1)-C(2)	109.0 ± 0.2	109.4 ± 0.5
N-C(2)-C(1)	110.1 ± 0.2	109.4 ± 0.7
N-C(2)-C(3)	113.5 ± 0.2	112.8 ± 0.6
C(1)-C(2)-C(3)	107.4 ± 0.2	105.4 ± 0.6
C(2)-C(3)-C(4)	105.5 ± 0.2	106.8 ± 0.6
S-C(4)-C(3)	105.8 ± 0.2	102.4 ± 0.6

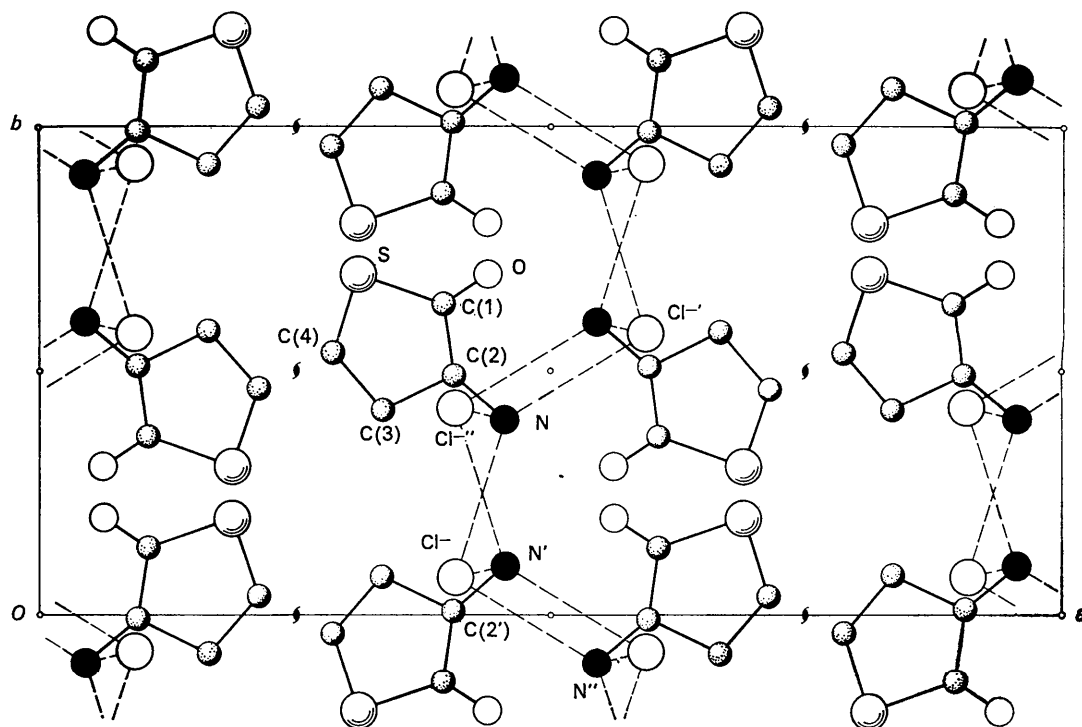
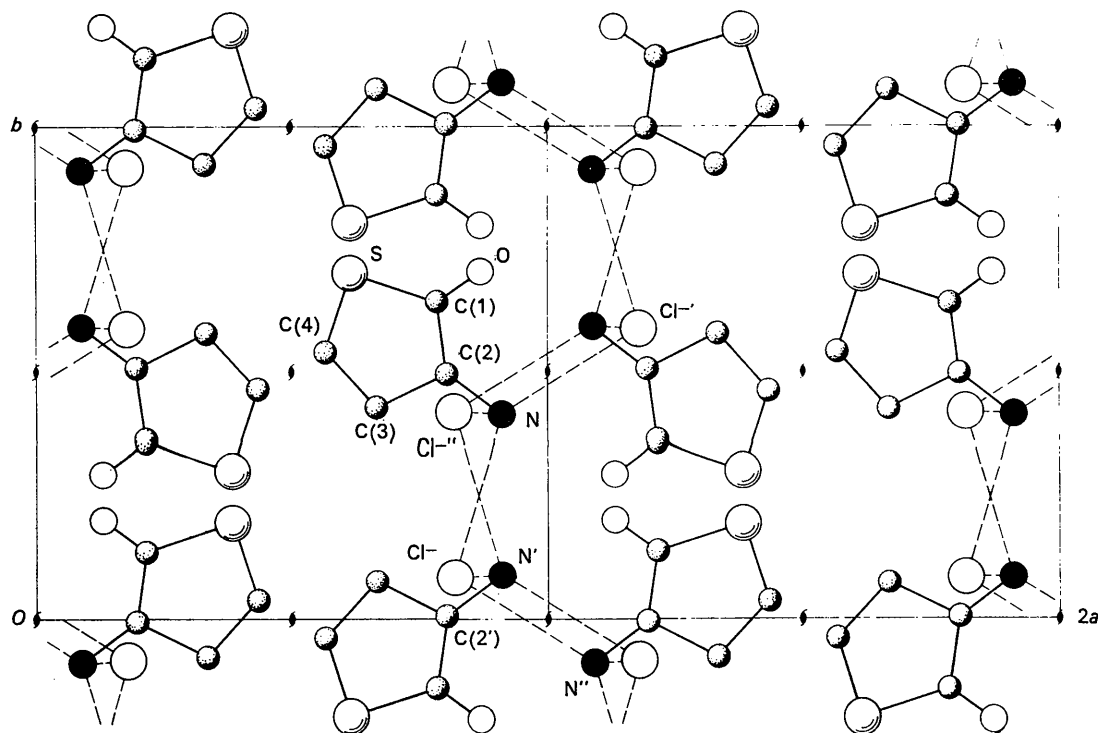
Fig. 2. Orthographic projection along the c axis of HTL- $Pbca$.Fig. 3. Orthographic projection along the c axis of HTL- $Pbc2_1$.

Table 8. Covalent bond lengths and angles involving hydrogen atoms and their estimated standard deviations for HTL-Pbca

Bond	Length
H(1)-C(2)	0.93 ± 0.03 Å
H(2)-C(3)	1.00 ± 0.03
H(3)-C(3)	1.00 ± 0.03
H(4)-C(4)	1.02 ± 0.03
H(5)-C(4)	1.02 ± 0.03
H(6)-N	0.98 ± 0.03
H(7)-N	1.00 ± 0.03
H(8)-N	0.88 ± 0.04
Angle	Degrees
H(1)-C(2)-C(1)	104.9 ± 1.6°
H(1)-C(2)-N	113.3 ± 1.5
H(1)-C(2)-C(3)	107.2 ± 1.5
H(2)-C(3)-C(2)	116.0 ± 1.5
H(2)-C(3)-C(4)	112.2 ± 1.6
H(2)-C(3)-H(3)	99.8 ± 2.2
H(3)-C(3)-C(2)	113.9 ± 1.4
H(3)-C(3)-C(4)	109.5 ± 1.4
H(4)-C(4)-C(3)	112.8 ± 1.4
H(4)-C(4)-S	106.7 ± 1.5
H(4)-C(4)-H(5)	115.6 ± 2.1
H(5)-C(4)-C(3)	110.3 ± 1.6
H(5)-C(4)-S	104.8 ± 1.6
H(6)-N-C(2)	110.2 ± 1.8
H(6)-N-H(7)	98.4 ± 2.3
H(6)-N-H(8)	111.9 ± 2.9
H(7)-N-C(2)	113.9 ± 1.6
H(7)-N-H(8)	105.3 ± 2.7
H(8)-N-C(2)	115.8 ± 2.1

Table 9. Hydrogen-bond lengths and their estimated standard deviations

Lengths involving hydrogen atoms of HTL-Pbc₂₁ are omitted because of the uncertainty of these positional parameters.

Bond	HTL-Pbca	HTL-Pbc ₂ ₁
N-Cl-	3.222 ± 0.003 Å	3.233 ± 0.007 Å
N-Cl'-	3.178 ± 0.002	3.175 ± 0.006
N-Cl''	3.185 ± 0.003	3.151 ± 0.008
Cl-H(6)	2.36 ± 0.03	—
Cl'-H(7)	2.22 ± 0.03	—
Cl''-H(8)	2.34 ± 0.04	—
Cl-C(2')	3.558 ± 0.003	3.640 ± 0.010
Cl-H(1')	2.77 ± 0.03	—

Table 10. Hydrogen-bond angles and their estimated standard deviations

Angles involving HTL-Pbc₂₁ hydrogen atoms are omitted because of the uncertainty of these positional parameters.

Angle	HTL-Pbca	HTL-Pbc ₂ ₁
N-Cl'-N'	91.5 ± 0.1	95.5 ± 0.2
N-Cl'-N''	104.8 ± 0.1	102.1 ± 0.2
N'-Cl'-N''	69.6 ± 0.1	88.0 ± 0.2
Cl'-N-Cl'-	134.8 ± 0.1	132.1 ± 0.2
Cl'-N-Cl''	98.8 ± 0.1	97.3 ± 0.2
Cl''-N-Cl''	110.4 ± 0.1	114.8 ± 0.2
H(6)-N-Cl-	23.1 ± 1.7	—
H(7)-N-Cl'-	14.7 ± 1.5	—
H(8)-N-Cl''	12.7 ± 2.1	—
H(1')-C(2')-Cl-	27.6 ± 1.5	—

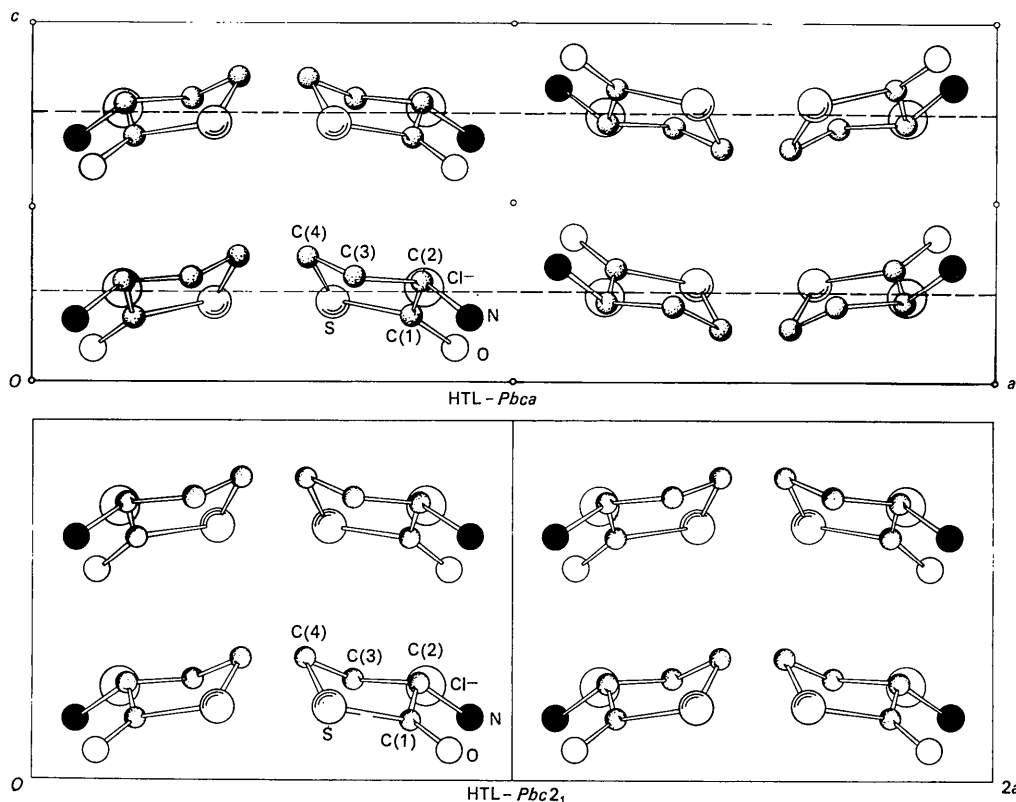


Fig. 4. Orthographic projections along the *b* axis of HTL-Pbca and HTL-Pbc₂₁.

in a plane and the fifth atom, C(3), displaced approximately 0.7 Å from this plane. For both HTL-*Pbca* and HTL-*Pbc*₂₁ the r.m.s. deviation of S, O, C(1), C(2), and C(4) from their least-squares plane is 0.02 Å; the maximum deviation is 0.026 Å. The thiolactone group exhibits the same characteristic planarity, and the associated shortening of the bond between the hetero atom and the C=O group, as has been found for lactones (McConnell, Mathieson & Schoenborn, 1964). Bond lengths and angles for HTL are similar to those found in other amino acids.

Molecular packing of the two crystal forms is very similar. Indeed, the [*c*] projections of the *Pbca* and *Pbc*₂₁ structures are almost indistinguishable (Figs. 2 and 3). In both polymorphs the molecules lie in layers at $z = \frac{1}{4}$ and $\frac{3}{4}$, with the plane of the thiolactone ring approximately normal to the unit-cell *c* axis. The principal packing difference is apparent in Fig. 4: the O, N ends of adjacent molecules are staggered on opposite sides of the a glide plane in HTL-*Pbca*, while they are directly opposed in HTL-*Pbc*₂₁. This produces slightly tighter packing in the *Pbca* polymorph, accounting for its slightly greater density and perhaps for its lower mean temperature factor (mean $B = 2.79$ Å² for HTL-*Pbca* and 4.45 Å² for HTL-*Pbc*₂₁).

The prominent feature of the molecular arrangement in both polymorphic crystal structures is an extensive hydrogen-bonding network between the ammonium groups and chloride ions. The three protons of the ammonium group participate in hydrogen-bond formation with three chloride ions. Since there is but one chloride ion in the asymmetric unit, each chloride ion must accept hydrogen bonds from three NH₃⁺ donor groups. In both polymorphs the chloride ion and its three hydrogen-bonded ammonium groups form a distorted triangular pyramid. One clear difference between the two structures is that the pyramid is more regular in HTL-*Pbc*₂₁, the three N-Cl-N angles at the Cl⁻ vertex being more nearly equal in this structure. This may be seen in the first three entries of Table 10, although it is not evident in the projections shown in Figs. 2, 3 and 4.

There is possibly a weak hydrogen bond between the chloride ion and C(2') in HTL-*Pbca* (atom designations correspond to those in Fig. 2). The Cl⁻-H(1') distance is approximately 0.2 Å shorter than the normal van der Waals contact of 3.0 Å. However, the Cl⁻-C(2')-H(1') angle of 27.6° is rather large and places the ex-

istence of such a hydrogen bond only on the verge of credibility.

Thermal ellipsoids were calculated for both polymorphs but are not listed here because of their dubious accuracy. It is probable that the thermal parameters were strongly affected by neglect of absorption corrections. All atoms, however, had significant anisotropy according to the criteria proposed by Kraut & Jensen (1963). In general the direction of greatest vibration for the ring atoms was approximately normal to the plane of the ring. Bond lengths were not corrected for librational motions of the entire molecule, since molecular libration is evidently not an important mode of thermal motion.

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